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MOLECULAR-DYNAMICS INVESTIGATION OF NONLINEAR DIELECTRIC RESPONSE FOR A FERROELECTRIC MODEL

By

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Abstract

The relation between the applied electric field E and its response P of a unified classical oscillator model for order-disorder and displacive ferroelectrics is investigated numerically on computer. It is shown that the relation between the E and the imaginary part of P depends on the type of the ferroelectrics. On the other hand, there are few differences between the two types in the static field case of the E - P relation.

Recently, Onodera¹⁾ has calculated the dynamic susceptibility of a unified oscillator model for order-disorder and displacive ferroelectrics with the aid of the linear-response theory.^{2),3)} He considered an assembly of classical oscillators moving in the anharmonic potential $V(x) = Ax^4 + Bx^2$, where x stands for the displacement of an oscillator. A is taken to be definitely positive, while B may be either positive or negative. The potential has one or two minima, depending on the sign of B . An interaction between these anharmonic oscillators is bilinear in their displacements and the interaction is treated in the Weiss approximation. The applied electric field E interacts with the dipole moment of the oscillator, which he assumes to be proportional to its displacement.

In the present letter, we shall numerically investigate the nonlinear dielectric response of a ferroelectric model which is similar to Onodera's model. Our model can be described by the Hamiltonian.

$$H = \sum_{i=1}^N \left\{ \frac{1}{2} M \dot{x}_i^2 + (Ax_i^4 + Bx_i^2) + \gamma \left(\frac{1}{N-1} \sum_{j \neq i}^N x_j \right) x_i + E \cos(\omega t) x_i \right\} + H', \quad (1)$$

Where N is the number of the oscillators, M is the mass of the oscillator, γ is the coupling constant and ω is the angular frequency of the applied electric field. H' is the coupling Hamiltonian between the system and a heat bath which is expressed by an ideal gas. We take the mass of the gas's particle 0.1 times as heavy as the mass of the

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oscillator and the particle of the gas collides with the oscillator. The collision produces a friction of the motion of the oscillator. The friction is a loss mechanism of our model which is not contained in Onodera's model. If we take the thermodynamic limit, namely N tend to infinity, our model undergoes a second-order phase transition in both $B > 0$ and $B < 0$ cases at a Curie temperature determined by γ^1 .

We define the electric polarization P as $P = (\sum_{i=1}^N x_i) / N$. The time dependence of P can be expressed as

$$P(t) = P' \cos(\omega t) + P'' \sin(\omega t), \quad (2)$$

where we neglect the higher harmonic terms.

According to Onodera we use the following units:

$$\text{Temperature} \dots\dots B^2 / 4Ak_B,$$

$$\text{Frequency} \dots\dots (2|B|/M)^{1/2},$$

where k_B is the Boltzmann constant.

We investigate the time evolution of the system by solving the equations of motion of the Hamiltonian ($H-H'$) numerically on computer by means of the Runge-Kutta method. The collision processes are inserted in the above time evolution about 10 times in a period of the oscillation in which we use random numbers. The random number expresses the velocity of the particle v of the one-dimensional ideal gas. Making use of the equipartition law of energy and $\bar{v} = 0$, the standard deviation of the velocity is determined by the temperature of the heat bath T and the mass of the particle m as $\{(\overline{v-\bar{v}})^2\}^{1/2} = \{\overline{v^2}\}^{1/2} = \{k_B T / m\}^{1/2}$. We calculate P' and P'' as a long time average, therefore, our results do not depend on the initial condition of our model. Our computer simulations show that the dependence of N is so small that we take $N=40$.

The polarization P versus the applied electric field E in displacive case, namely $B > 0$, at several temperatures with $\omega=0$ is shown in Fig. 1. The same relation of the order-disorder case, namely $B < 0$, is shown in Fig. 2. Fig. 1 and 2 show that there are few differences between the order-disorder case and the displacive case in the static nonlinear dielectric response. The E - P relation with $\omega=0.25$ of the displacive case and order-disorder case are shown in Fig. 3 and 4, respectively. In both cases, the temperatures are taken near their Curie points which are estimated by Fig. 1 and 2. It is interesting that in dynamic case, the relation E - P'' depends on the type of the ferroelectrics.

Onodera's model becomes non-ergodic when $B < 0$.¹⁾ It is a difficult question whether we can apply the linear-response theory to the non-ergodic case or not. In our case, however the coupling Hamiltonian H' enables us to remove the problem of the non-ergodicity. In usual cases, the motion of the oscillation accompanies a friction which gives a loss. Due to the convenience of the calculation, we express the friction by H' .

Further studies including the friction dependence of the response are now going

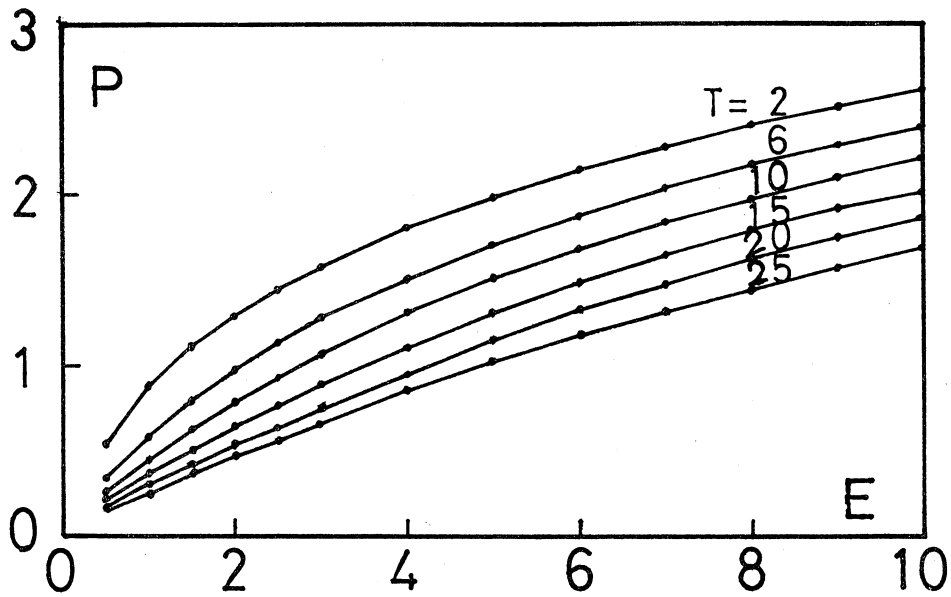


Fig. 1. The static E - P relation of the displacive case with $A=1$, $B=2$, $\gamma=-2$ and $N=40$ at several temperatures.

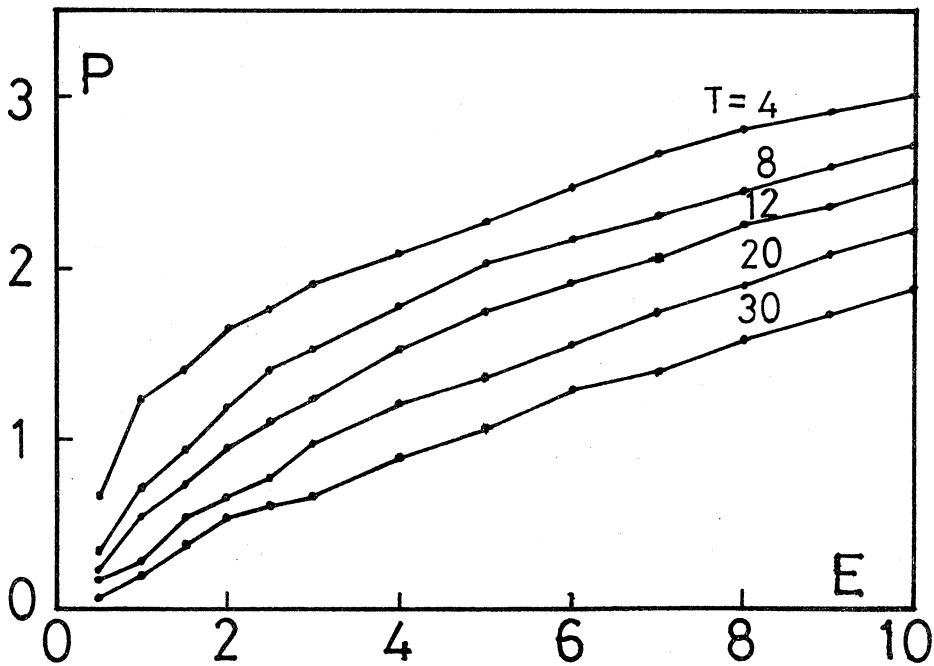


Fig. 2. The static E - P relation of the order-disorder case with $A=1$, $B=-2$, $\gamma=-0.1$ and $N=40$ at several temperatures.

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The numerical calculations were performed using a FACOM 230-75 computer at the Kyushu University Computing Center and a FACOM 230-45S computer at the Kagoshima University Computing Center.

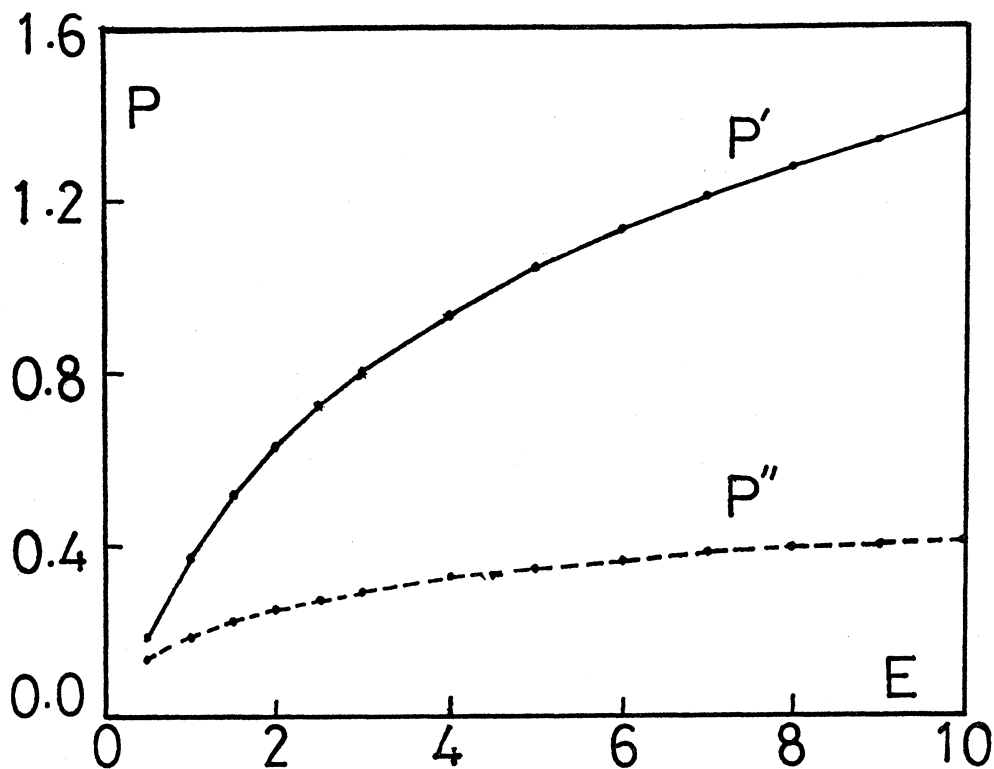


Fig. 3. The E - P' and E - P'' relations of the displacive case with $A=1$, $B=2$, $\gamma=-2$, $N=40$ and $\omega=0.25$ at $T=2$.

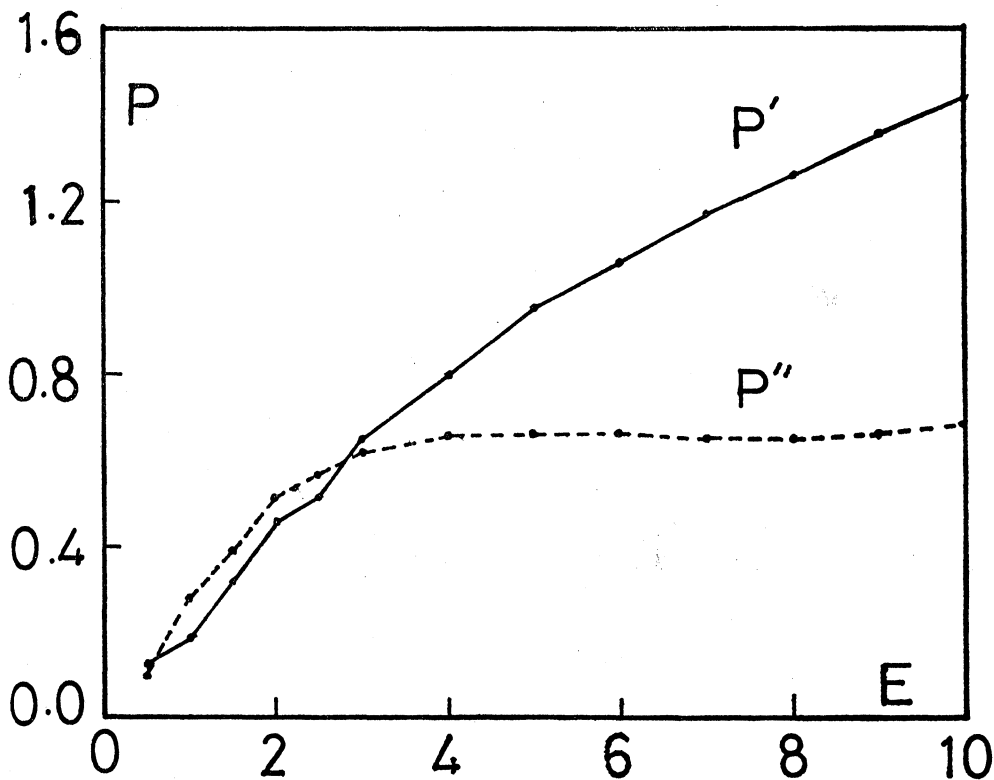


Fig. 4. The E - P' and E - P'' relations of the order-disorder case with $A=1$, $B=-2$, $\gamma=-0.1$, $N=40$ and $\omega=0.25$ at $T=4$.

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